

Time-Resolved Fluorescence and Essential Dynamics Study on the Structural Heterogeneity of p53DBD Bound to the Anticancer p28 Peptide

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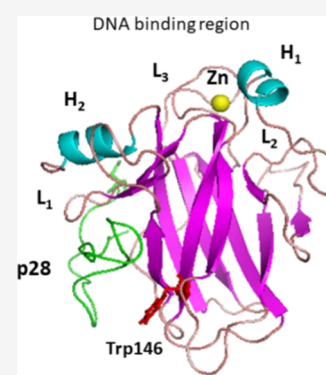
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ABSTRACT: Time-resolved fluorescence emission was combined with molecular dynamics (MD) simulations to investigate the DNA-binding domain (DBD) of the tumor suppressor p53 alone and its complex with the anticancer peptide p28 (DBD/p28). The fluorescence emission decay of the lone Trp residue, from both DBD and DBD/p28, was well-described by a stretched exponential function. Such a behavior was ascribed to heterogeneity in the Trp relaxation behavior, likely due to the coexistence of different conformational states. The increase of the stretching parameter, on passing from DBD to DBD/p28, indicates a reduced heterogeneity in the Trp146 environment for DBD/p28. Moreover, the effects of p28 on the global dynamics of DBD were analyzed by the essential dynamics method on 30 ns long MD trajectories of both DBD and DBD/p28. We found the establishment of wide-amplitude anharmonic modes throughout the DBD molecule, with a particularly high amplitude being detected in the DNA-binding region. These modes are significantly reduced when DBD is bound to p28, consistently with a structure stabilization. In summary, the results indicate that p28 binding has a strong effect on both the local and global heterogeneity of DBD, thus providing some hints to the understanding of its anticancer activity.



INTRODUCTION

It is now well ascertained that the structure–function paradigm, central in the biology of proteins, is to be extended to include intrinsically disordered proteins (IDPs), which possess (a few or many) intrinsically disordered regions.¹ IDPs are present in living organisms with a higher frequency as far as the complexity level of organisms increases.² The partial lack of a well-defined structure confers to these proteins some advantages in terms of functionality. Indeed, thanks to the coexistence in solution of several widely different conformations, IDPs acquire the capability to bind different biological molecules (proteins, DNA, and RNA), by constituting then a sort of hub with a crucial role in many biological processes.^{2,3}

The tumor suppressor p53 is a very important IDP and is at the center of a complex interaction network regulating numerous cellular processes, such as cell cycle progression, apoptosis induction, and DNA repair.^{4,5} Under physiological conditions, p53 possesses both structured and disordered domains.⁶ More specifically, p53 is a tetrameric protein composed of four identical subunits, with each monomer consisting of an N-terminal transactivation domain (NTD), a C-terminal domain (CTD), and a core domain, with large structured regions, which binds DNA (DBD).^{7–10} Since the oncosuppressive function of p53 is reduced or inactivated in many human cancers, large efforts are devoted to identify, design, and potentiate molecules able to stabilize or restore the anticancer activity of p53, with particular attention to its DBD,

which is incidentally mutated in more than half of the human cancers.^{11–13}

We have investigated the interaction between full-length p53, or its isolated domains (DBD or even NTD), and the blue-copper-containing protein Azurin (from *Pseudomonas aeruginosa* bacterium), which has been found to preferentially enter cancer cells and form a specific and stable complex with p53, with a concomitant increase of its tumor suppressor activity.^{14–19} More recently, we have focused our attention on the p28 peptide, which is formed by the α -helix fragment (50–77 AA) of Azurin and is endowed with an efficient cellular penetration ability. Upon entering the cell, p28 can bind to p53, leading to an increase of its intracellular levels. Indeed, p28 inherits the overall tumoricidal activity of Azurin but with lesser side effects.²⁰ The interaction between the p28 peptide and DBD has been investigated by atomic force spectroscopy (AFS), surface plasmon resonance (SPR), Raman, fluorescence, and Förster resonance energy transfer (FRET) techniques; in the last case, DBD containing a lone tryptophan residue (Trp146) has been used. These studies have allowed us

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66 to demonstrate the formation of a stable complex between
67 DBD and p28 and evaluate the corresponding binding affinity.
68 Additionally, the distance between the donor (the lone Trp146
69 residue) and the acceptor (a suitable dye bound to the N-
70 terminal of p28) has been estimated. Knowledge of this
71 distance combined with the results from a docking study,
72 refined by molecular dynamics (MD) simulations, has allowed
73 us to single out the best model for the DBD-p28 complex.²¹

74 In this context, it would be interesting to investigate the
75 structural and dynamical effects that are induced by the
76 binding of p28 to DBD. To address these aspects, we have
77 applied time-resolved fluorescence combined with molecular
78 dynamics (MD) simulations on both DBD and its complex
79 with the p28 peptide (DBD/p28). We found that the time-
80 resolved fluorescence emission of the lone Trp146 of DBD is
81 characterized by a stretched exponential decay, indicative of a
82 heterogeneous relaxation mechanism of the Trp residue. Such
83 a result has been put into relationship to the coexistence, in
84 solution, of different conformational states, which can be
85 ascribed to the partial disorder character of DBD. The
86 conformational heterogeneity around Trp146 has been
87 observed to persist in DBD bound to p28, but to a lower
88 extent, with this being indicative of a fine tuning of the
89 structural disorder of DBD as induced by binding of p28. To
90 elucidate the p28 effects on the global dynamics of DBD, we
91 have carried out MD simulations on DBD and DBD/p28,
92 within the same temporal window of the fluorescence
93 experiments. Essential dynamics (ED) analysis of the
94 corresponding MD trajectories has put into evidence the
95 existence of wide anharmonic modes within the DBD
96 molecule. The amplitudes of these modes undergo a significant
97 reduction when DBD is bound to p28. All these results,
98 consistent with the intrinsically disordered character of the
99 biomolecule, could be of some help for the biophysical
100 interpretation of the anticancer role of p28, as exerted upon its
101 interaction with DBD.

102 ■ EXPERIMENTAL AND COMPUTATIONAL 103 METHODS

104 **Materials.** Recombinant human DNA-binding domain
105 (DBD) of p53 (residues 94–300) (23400 Da, purity > 90%,
106 verified by high-performance liquid chromatography (HPLC)
107 and MS by the producer), containing a lone tryptophan
108 (Trp146), was purchased from Genscript (Piscataway, NJ)
109 using the BacPower Guaranteed Bacterial Protein Expression
110 Service.

111 The p28 peptide (LSTAADMQGVVTDGMASGLDK-
112 DYLPDD, 2914 Da) was synthesized and verified by MS
113 and HPLC (93% purity) by GenScript (Piscataway, NJ).
114 Human serum albumin (HSA) from human serum lyophilized
115 powder, essentially globulin free with a purity degree higher
116 than 99%, was purchased from Sigma Aldrich and used without
117 further purification. Phosphate-buffered saline (PBS, 50mM at
118 pH 7.4) was used as buffer.

119 **Absorbance, Static, and Time-Resolved Fluorescence**
120 **Measurements.** Absorbance spectra were recorded at room
121 temperature by a double-beam Jasco V-550UV/visible
122 spectrophotometer using 1 cm path length cuvettes and 1
123 nm bandwidth in the spectral region 220–750 nm; PBS was
124 used as reference.

125 Steady-state fluorescence measurements were performed
126 with a FluoroMax-4 spectrofluorometer (Horiba Scientific,
127 Jobin Yvon, France). The samples were excited at 295 nm and

fluorescence emission was collected from 305 to 580 nm using
128 1 nm increments and an integration time of 0.50 s. A 2 nm
129 bandpass was used in both excitation and emission paths. 130
Spectra were acquired in the signal to reference (S/R) mode to
131 take into account for random lamp intensity fluctuations. 132
Emission spectra were corrected for Raman contribution from
133 the PBS. 134

Time-resolved fluorescence measurements were performed
135 at room temperature with the time-correlated single photon
136 counting method using the abovementioned FluoroMax-4
137 spectrofluorometer, operating at a repetition rate of 1 MHz
138 and running in reverse mode. The apparatus was equipped
139 with a pulsed nanosecond light-emitting diode (LED) 140
excitation head (Horiba Scientific, Jobin Yvon), emitting at
141 295 nm with a bandwidth of 20 nm (which was lowered to 4
142 nm by the slits) and a temporal width lower than 1 ns. 143
Fluorescence lifetime data were detected at 345 nm and
144 acquired until the peak signal reached 10k counts and were
145 analyzed using the impulse response function (DAS6 software,
146 Horiba Scientific). For each fluorescence experiment, 10
147 measurements were performed on independently prepared
148 samples by determining the average and the corresponding
149 standard deviation. 150

Analysis of Time-Resolved Fluorescence. Fluorescence
151 decays were preliminarily analyzed by the sum of exponential
152 components by employing a non-linear least square analysis
153 including deconvolution of the prompt from Horiba, Jobin
154 Yvon 155

$$156 I(t) = a_0 + \sum_{i=1}^n a_i e^{(-t/\tau_i)} \quad (1)$$

where $I(t)$ is the time-dependent intensity, a_0 gives the
157 background, and a_i are pre-exponential factors representing
158 fractional contribution to the time-resolved decay of the i th
159 component with lifetime τ_i . The goodness of the fit was judged
160 in terms of both the χ^2 value and weighted residuals. The
161 average fluorescence lifetime $\langle \tau \rangle$ was then calculated by
162

$$163 \langle \tau \rangle = \frac{\sum_{i=1}^n a_i \tau_i}{\sum_{i=1}^n a_i} \quad (2)$$

Modeling Procedures. The structure of DBD was derived
164 from the B chain of DBD in complex with a consensus DNA (1
165 TUP entry from the protein data bank).¹⁰ The structure was
166 suitably adjusted to match the DBD portion used in
167 fluorescence and FRET experiments. In particular, the short
168 290–300 AA portion was added using the I-TASSER suite.²² 169
The structure of DBD is composed of two antiparallel β -sheets
170 of four and five strands, respectively, forming a β -sandwich
171 structure (see Figure 1). The Zn ion is tetrahedrally
172 coordinated to the side chains of Cys176, His179, Cys238,
173 and Cys242, forming a Zn-finger motif, which is connected to
174 the L₂ and L₃ loops.²³ The interaction of the Zn ion with its
175 ligands was treated through a bonded approach, in which the
176 Zn–N and Zn–S bonds and S–Zn–S angles were set
177 according to the parameters provided in refs 21, 24–26. The
178 binding of DBD to DNA occurs within the H₂ and the L₁ and
179 L₃ loops in a region, conventionally chosen to be the northern
180 part of the molecule. 181

The initial coordinates of the DBD-p28 complex were taken
182 from the “best model” as obtained by a previous study based
183 on a computational docking approach exploiting the structural
184

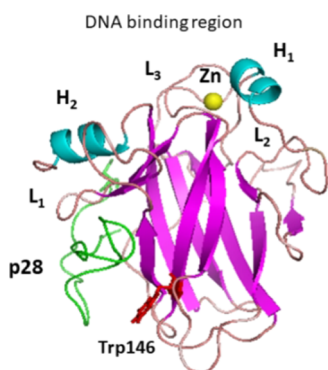


Figure 1. Structure of the complex between DBD (magenta) and p28 (green), according to the “best model” obtained by the method reported in ref21; the Zn ion is represented by a yellow sphere and Trp146 in blue.

185 information derived from FRET data,²¹ with such a model
186 being shown in Figure 1.

187 Molecular Dynamics (MD) Simulations and Analysis.

188 MD simulations of DBD and the DBD-p28 complex in water
189 were carried out by the GROMACS 2018 package²⁷ using
190 AMBER03 Force Field DBD and p28²⁸ and SPC/E for
191 water.²⁹ DBD and DBD-p28 were centered in a cubic box of
192 edge 9.0 nm. Simulations were performed by following the
193 procedures described in refs 21, 30. Briefly, the boxes were
194 filled with water molecules to reach a hydration level of 9 g
195 water/g protein. The ionization states of protein residues were
196 fixed at pH 7, and Cl⁻ ions (one for DBD and five for DBD/
197 p28) were added to keep the systems electrically neutral. H
198 bonds were constrained with the LINCS algorithm,³¹ while the
199 particle mesh Ewald (PME) method^{32,33} was used to calculate
200 the electrostatic interactions with a lattice constant of 0.12 nm.
201 Periodic boundary conditions in the NPT ensemble with $T =$
202 300 K and $p = 1$ bar with a time step of 1 fs were used. The
203 temperature was controlled by the Nosé–Hoover thermostat
204 with a coupling time constant τ_T of 0.1 ps,³⁴ while Parrinello–
205 Rahman extended ensemble, with a time constant t_p of 2.0 ps,
206 was used to control the pressure.³⁵ Each system was minimized
207 and then heated to 300 K with steps at 50, 100, 150, and 250
208 K. MD trajectories were analyzed by the GROMACS package
209 tools.²⁷ Each model was submitted to a 30 ns long MD
210 trajectory, replicated 10 times. The temporal evolution of the
211 systems was monitored by following the root mean square
212 displacement (RMSD), root mean square fluctuations
213 (RMSF), surface accessible surface (SAS) area, and the

214 methods described in ref36. The equilibration of each system
215 was checked by analyzing the RMSD as a function of time. The
216 figures were created with Pymol and VMD.^{37,38}

Essential Dynamics (ED) Analysis. The sampled MD
217 trajectories were analyzed by ED, a method based on principal
218 component analysis (PCA) aimed at identifying a new
219 reference frame to describe the overall dynamics of the system
220 and then allowing to extract the protein motions mostly
221 contributing to the overall dynamics.³⁹ ED was carried out by
222 the Covar and Ana eig subroutines of GROMACS.^{40,41} After a
223 least-square fit to remove roto-translations, the covariance
224 matrix C_{ij} was calculated by
225

$$C_{ij} = \langle (x_i - \langle x_i \rangle)(x_j - \langle x_j \rangle) \rangle \quad (3)$$

227 where x_i are the x , y , and z Cartesian coordinates of the C_α
228 atoms of DBD and $\langle \rangle$ indicates the average over the
229 trajectories. C_{ij} was diagonalized by finding out a set of
230 eigenvalues and eigenvectors; the eigenvectors correspond to
231 directions in an N -dimensional space and are called principal
232 coordinates (PCs), while the eigenvalues represent the total
233 mean square fluctuation along the corresponding eigenvectors.

234 RESULTS AND DISCUSSION

Fluorescence Emission. Figure 2A shows the emission
235 spectrum of DBD (black line) and DBD after the addition of
236 p28 (DBD/p28; red line) in a 1:1 molar ratio. Both spectra
237 have been obtained by excitation at 295 nm, at which the lone
238 tryptophan residue (Trp146) from DBD still absorbs, while the
239 Tyr and Phe residues, from DBD and p28, are not substantially
240 excited.
241

242 The spectrum of DBD is peaked at about 345 nm (see the
243 arrow), pointing out that Trp146 is almost fully exposed to the
244 solvent, this being consistent with the SAS for Trp146 of about
245 0.8 nm², as evaluated from the DBD X-ray structure.⁴² The
246 fluorescence emission intensity of DBD/p28 is reduced with
247 respect to that of DBD, although no significant wavelength
248 shift of the peak is observed. In particular, we found (32 ± 2)
249 $\times 10^3$ counts for DBD and $(25 \pm 2) \times 10^3$ counts for DBD/
250 p28, by averaging the results over 10 samples. Such a 20%
251 decrease in the fluorescence intensity indicates that p28 affects
252 the fluorescence emission of DBD without changing the
253 solvent exposition of Trp146. This result finds a correspond-
254 ence with the formation of a stable complex between DBD
255 p28, in which Trp146 is not directly involved, as previously
256 found.²¹ Indeed, we recall that static fluorescence experiments
257 demonstrated the formation of a stable complex between DBD

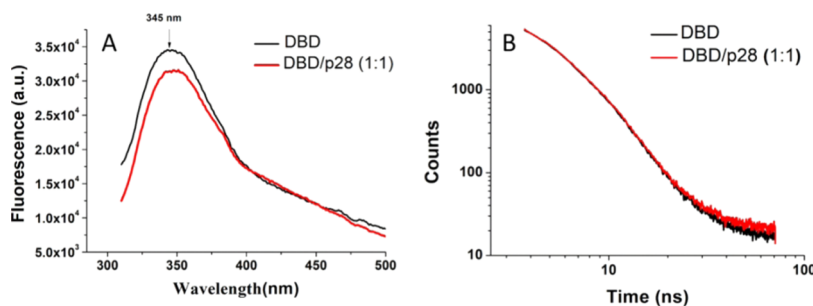


Figure 2. (A) Fluorescence emission spectra of DBD (black line) and DBD in the presence of p28 (1:1) (red line), registered upon excitation at 295 nm, and corrected for the Raman scattering of the buffer. (B) Time-resolved fluorescence emission of DBD (black line) and DBD in the presence of p28 (1:1) (red line). The decays have been obtained by exciting the samples with a 295 nm nanoLED pulse, while emission was collected at 345 nm.

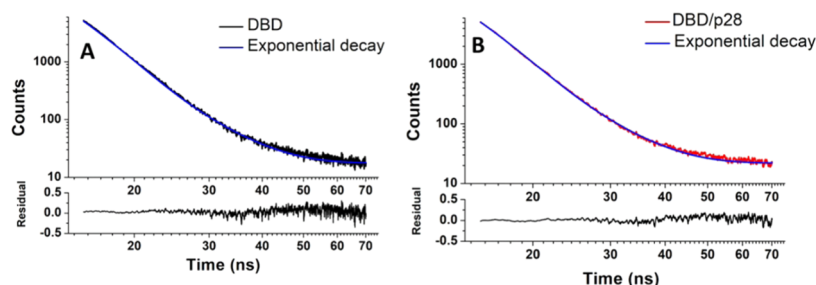


Figure 3. Fits (blue curves) by the sum of two exponential decays through eq 1 of fluorescence emission data for (A) DBD (black curve) and (B) DBD/p28 (red curve); the fitting parameters are reported in Table 1.

Table 1. Fitting Parameters of Time-Resolved Fluorescence Emission Data for DBD and DBD/p28, as Obtained by Averaging the Results over 10 Curves, by the Sum of Two Exponentials through eq 1

sample	τ_1 (ns)	τ_2 (ns)	$\langle\tau\rangle$ (ns) (eq 2)	a_1	a_2	χ^2 range
DBD	2.63 ± 0.03	7.32 ± 0.03	3.24 ± 0.03	36.3 ± 0.1	5.5 ± 0.1	1.1–1.2
DBD/p28	2.63 ± 0.03	7.34 ± 0.03	3.21 ± 0.03	35.5 ± 0.1	5.3 ± 0.1	1.1–1.2

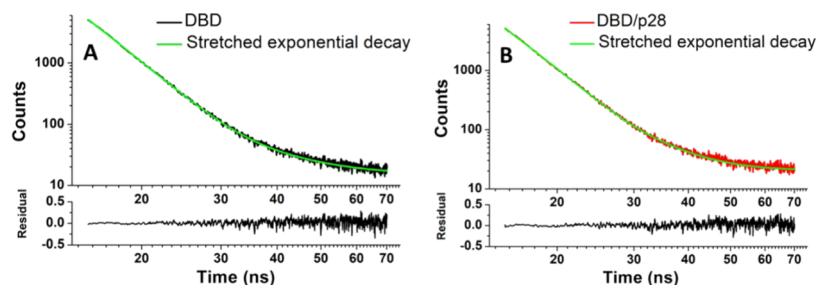


Figure 4. Fits (green curves) by the sum of a simple exponential decay and a stretched exponential decay through eq 4 of fluorescence emission data for: (A) DBD (black curve) and (B) DBD/p28 (red curve); the fitting parameters are reported in Table 1.

Table 2. Fitting Parameters of Time-Resolved Fluorescence Emission Data for DBD and DBD/p28, as Obtained by Averaging the Results from 10 Curves, by the Sum of an Exponential and a Stretched Exponential through eq 4

sample	τ_1 (ns)	τ_2 (ns)	β	a_1	a_2	χ^2 range
DBD	2.63 ± 0.04	3.56 ± 0.03	0.74 ± 0.01	2.4 ± 0.1	2.4 ± 0.1	1.0–1.1
DBD/p28	2.63 ± 0.03	3.55 ± 0.03	0.82 ± 0.01	2.4 ± 0.1	2.4 ± 0.1	1.0–1.1

258 and p28 with a dissociation constant K_D of $\sim 7.4 \times 10^{-6}$ M.²¹
 259 Figure 2B representatively shows the time-resolved fluores-
 260 cence emissions of DBD (black line) and DBD/p28 (red line),
 261 excited at 295 nm by a nanosecond pulsed LED. The two
 262 fluorescence decays appear practically identical at times shorter
 263 than 30 ns, while they differ more than 10% at longer times.
 264 At first, the emission data of both DBD and DBD/p28 have
 265 been fitted by the sum of two exponential decays using eq 1.
 266 Although a satisfactory description of the global decay for both
 267 the samples has been obtained for shorter decay times, a slight
 268 deviation between experimental and fitted curves is detected at
 269 longer times (see an example in Figure 3A,B). The fitting
 270 parameters, extracted by eq 1, and averaged from a collection
 271 of 10 independently prepared samples, are reported in Table 1.
 272 The resulting lifetimes are practically the same for both the
 273 samples, and they are in agreement with the values previously
 274 obtained ((3.19 ± 0.01) and (3.17 ± 0.02) ns), for DBD and
 275 DBD/p28, respectively.^{21,43} We also note that the amplitudes
 276 of the two pre-exponential factors are very similar for the two
 277 samples, thus indicating that the binding of p28 does not
 278 modify the decay populations in DBD.
 279 To more precisely describe the time-resolved fluorescence
 280 emission, and by taking into consideration the disordered

character of DBD, we were led to consider the possibility of a
 stretched exponential decay, commonly used to describe the
 fluorescence emission of heterogeneous systems, for both the
 fast and slow decay components.⁴⁴ We found that the
 fluorescence emission intensity of the DBD and DBD/p28
 samples can be well fitted by the sum of a simple exponential
 decay for the fast component and a stretched exponential
 decay for the slow one. This could be justified by the reported
 evidence that protein motions are characterized by an almost
 homogeneous behavior, while the slower one is almost
 heterogeneous, reflecting collective motions.⁴⁵ In particular,
 the following expression has been used

$$I(t) = a_0 + a_1 e^{(-t/\tau_1)} + a_2 e^{(-t/\tau_2)^\beta} \quad (4)$$

where a_0 describes the background, a_1 and a_2 are the pre-
 exponential factors for the simple and stretched exponential
 decays, respectively, τ_1 and τ_2 are the corresponding lifetimes,
 and β is the stretching parameter, which should fall between 0
 and 1, with the stretching character being progressively more
 relevant as far as it becomes less than 1.

Representative fits of the experimental time-resolved
 fluorescence emission using eq 4 are shown in Figure 4. At
 visual inspection, the agreement between fitted curves and

303 experimental data is very good throughout the time range and
304 it is significantly improved in comparison to that obtained
305 using two exponentials (see the residual plots in Figures 3 and
306 4) and witnessed by the lower chi-square values (see Tables 1
307 and 2). The fitting parameters, averaged from a collection of 10
308 independently prepared samples, are reported in Table 2. The
309 fast component, following a simple exponential decay, is
310 characterized by the same decay time for both the samples,
311 such a time being also equal to that emerging from the fit with
312 two exponential decays. Also, in this case, the amplitudes of the
313 two pre-exponential factors are substantially the same for the
314 two samples. The stretched exponential decay is characterized
315 by almost the same decay times for DBD and DBD/p28, with
316 the β parameter significantly deviating from 1 for both the
317 samples. However, β is slightly lower for DBD ($\beta = 0.74$)
318 compared to DBD/p28 ($\beta = 0.82$).

319 A stretched exponential decay, with a β parameter lower
320 than 1, generally reflects a glass-like character of the system,
321 and it has been observed in many disordered systems,
322 including biomolecules.⁴⁶ In proteins, a stretched exponential
323 function has been used to describe different processes, such as
324 photodissociation, protein, and hydration water dynamics, and
325 it has been put into relationship to the sampling of dynamical
326 states differing among them for structural details (usually
327 called conformational substates).^{45,47,48} The emission fluo-
328 rescence decay of Trp residues in proteins has been fitted by a
329 stretched exponential and interpreted in terms of a
330 heterogeneous environment around Trp residues, which may
331 include contributions from both the protein matrix and the
332 solvent.^{44,49,50} For single-tryptophan proteins, as in our case,
333 the observation of a stretched exponential for the emission
334 fluorescence decay of Trp could be directly related to the
335 behavior of a well-defined region of the protein rather than to
336 the complexity arising from the superposition of different
337 contributions. In this respect, for comparison, we have also
338 analyzed the time-resolved fluorescence emission decay of the
339 HSA protein, which contains a lone Trp (Trp214) exposed to
340 the solvent, with an SAS value of about 0.7 nm², as evaluated
341 from the X-ray structure.⁵¹

342 Figure 5 shows the time-resolved fluorescence emission data
343 of HSA (black curve) excited at 295 nm by a nanoLED pulse;

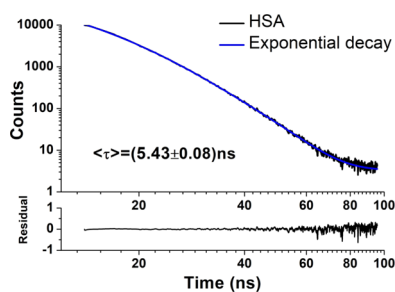


Figure 5. Fluorescence decay (black lines) obtained by exciting with a 295 nm nanoLED pulse; emission was collected at 345 nm. The blue line is a fit by the sum of two exponential decays. The average lifetime, as evaluated by eq 2, is reported.

344 the corresponding fit using the sum of two exponential decay
345 (red curve) is also shown. The fit almost perfectly describes the
346 data and the resulting average lifetime of (5.43 ± 0.08) ns,
347 which is in very good agreement with the literature (including
348 our previous results).^{52,53} The fit by the sum of exponentials
349 for HSA supports the hypothesis that the appearance of a

stretched exponential decay in DBD and DBD/p28 can be
ascribed to some heterogeneity in the Trp146 environment
arising from the protein matrix. In other words, the observation
of a stretched exponential decay in DBD can be put into
relationship to a local structural disorder of the protein, as
monitored by Trp146. On such a basis, the slight increase of
the β exponent observed for the Trp146 emission of DBD/p28
with respect to DBD indicates a decrease in the structural
disorder of DBD upon its binding to p28. In summary, the
time-resolved fluorescence emission decay of Trp146 in DBD,
alone and bound to p28, shows a heterogeneous relaxation
behavior, which reflects a local disorder undergoing a decrease
when DBD is bound to p28. Furthermore, since the p28
binding is not directly involved with Trp, as previously
demonstrated,²¹ these results suggest that p28 can affect the
global dynamics of DBD.

MD simulation results. To closely investigate the effects
induced by p28 on both the whole structure of DBD and the
Trp146 environment, we have performed MD simulations on
both DBD alone and the DBD/p28 complex. For each system,
we have collected MD simulation runs of 30 ns; such a time
interval is comparable with the temporal windows of the
fluorescence emission. For the DBD/p28 complex, whose
structure is not known, we have used the model previously
derived from FRET data (see the section Experimental and
Computational Methods and ref21).

First, we have assessed the stability of the complex by
monitoring the distance between DBD Trp146 and the p28
Leu19 residue, the latter one being the p28 residue closest to
DBD. Figure 6A shows the trend of this distance as a function
of time. For all of the runs, we noted that the distance varies in
time between 0.8 and 1.0 nm. These results indicate that p28
remains in the close proximity of DBD during the run, and
then they support the stability of the DBD/p28 complex.

Figure 6B shows a representative temporal evolution of the
RMSDs, averaged over all atoms, for DBD alone (black curve)
and for DBD/p28 (red curve). In both cases, the RMSDs
rapidly increase by reaching a plateau at about 0.25 nm within
2 ns, consistently with a quick relaxation of the structure.
Successively, we note an almost constant regime, in which fast
fluctuations are superimposed on slower ones. To avoid
instabilities due to protein equilibration, we have discarded the
first 3 ns, and the successive 27 ns long time intervals have
been then taken into consideration for further analysis.

Figure 7A,B shows the time-averaged RMSDs and RMSFs,
respectively, as a function of the residue number of DBD alone
(black lines) and of DBD/p28 (red lines), the secondary
structure being sketched at the bottom. RMSDs and RMSFs
exhibit similar profiles with almost all of the peaks located at
the same positions along the protein chain. Similar trends have
been observed for the other runs. The largest part of the peaks
appears in correspondence of the turns and bends, consistently
with the higher flexibility of these regions. The RMSDs in the
region involved in the binding with p28, marked with blue stars
in Figure 7, show slightly larger deviations from the initial
structure, consistent with a rearrangement of the DBD
structure upon binding the p28. At the same time, a reduction
in the RMSF values occurs, with this reflecting a decreased
mobility of residues in close contact with the peptide.
Additionally, for the region around Trp146 (marked with
circles in Figure 7), the RMSDs show both increase and
decrease on passing from DBD to DBD/p28, consistent with
slight structural changes of the Trp146 environment. At the

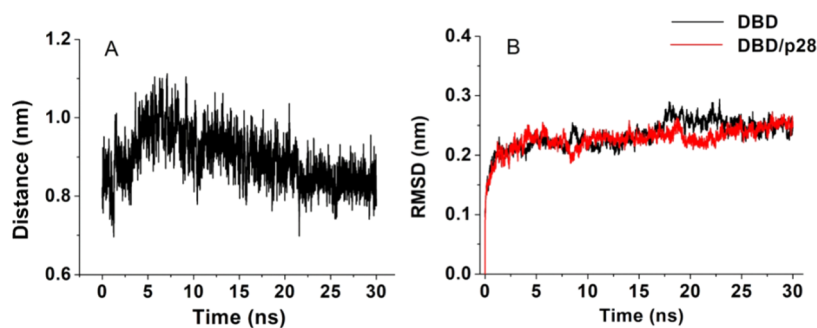


Figure 6. (A) Distance between Trp146 (DBD) and Leu19 (p28) as a function of time. (B) RMSD, averaged over all atoms, as a function of time for DBD alone (black line) and complexed with p28 (red line).

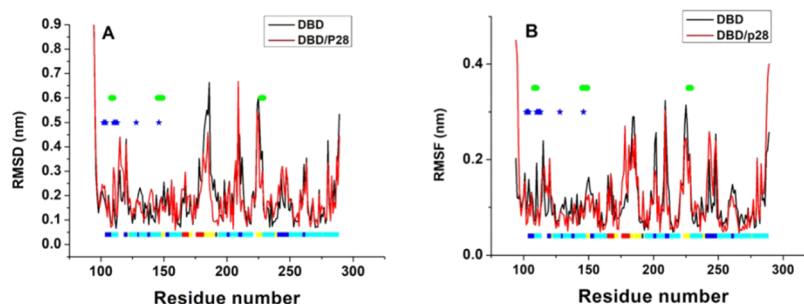


Figure 7. (A) RMSD and (B) RMSF as a function of the residue number for DBD (black lines) and DBD complexed with p28 (red lines). The secondary structure is sketched in the bottom: turn (blue squares), α -helices (red squares), bend (yellow squares), and β -sheets (cyan squares). Blue stars mark the residues involved in the binding with p28, while green circles mark the residues with a distance lower than 0.5 nm from Trp146.

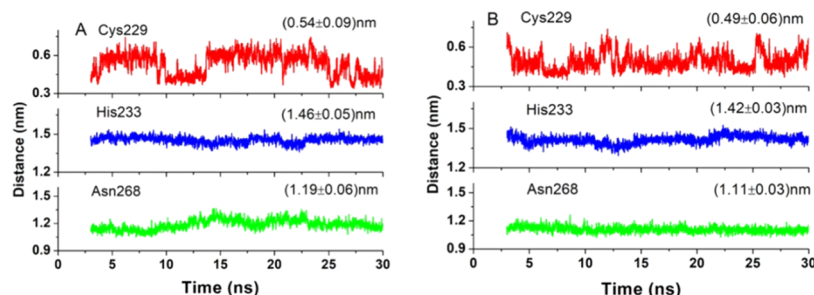


Figure 8. Temporal evolution of the distance between Trp146 and Cys229 (red curves), His233 (blue curves), and Asn268 (green curves) for DBD (A) and DBD/p28 (B).

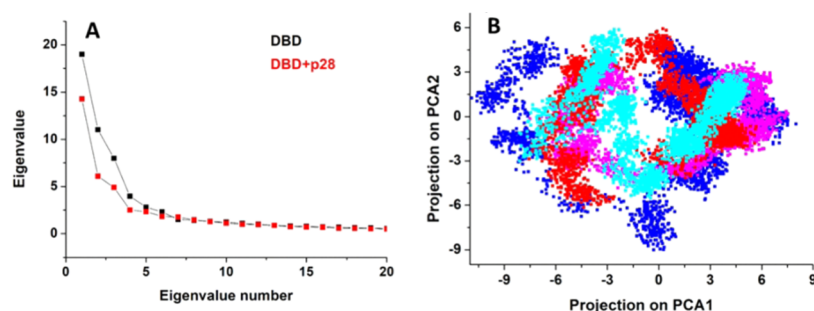


Figure 9. (A) Eigenvalues as a function of the eigenvalue number for DBD and DBD/p28. (B) 2D plot of the projection of PC1 versus the projection of PC2 for DBD (blue and magenta spots) and DBD/p28 (red and cyan spots).

413 same time, the RMSFs undergo a slight decrease, indicating a
 414 lower flexibility for this region when DBD is bound to p28.
 415 Then, we more closely analyzed the region around Trp146 to
 416 elucidate possible changes affecting the fluorescence properties
 417 of DBD. Accordingly, we monitored the distance between
 418 Trp146 and the Cys229, His233, and Asn268 residues, all of

them being located within 1.5 nm from Trp146 and assumed 419
 to be more susceptible to quench the Trp fluorescence.⁵⁴ 420
 Representative plots for the temporal evolution of these three 421
 distances are shown in Figure 8. We note fast fluctuations on 422
 which some jumps are superimposed; a higher number of 423
 jumps are detected for Cys229, which is the closest residue to 424

425 Trp146. The average distances and the standard deviations
426 (see Figure 8) are indicative that small changes occur upon
427 binding of p28 to DBD, with this being consistent with a
428 structural rearrangement of the protein region around Trp146.
429 Interestingly, the standard deviations of the analyzed distances
430 are found to be slightly lower on passing from DBD to DBD/
431 p28. Such a finding could be attributed to the reduction of the
432 structural heterogeneity around Trp146, which is in agreement
433 with what has been found by fluorescence.

434 To get insights into the global dynamics of DBD and DBD/
435 p28, we analyzed their MD trajectories by the ED method,
436 which restricts the conformational space, allowing us to find
437 out the relevant protein motions. The first 20 eigenvalues
438 extracted from the covariance matrix of the atomic fluctuations
439 and ranked in a decrement order, for DBD (black squares) and
440 DBD/p28 (red squares), are shown in Figure 9A. In both
441 cases, the eigenvalues rapidly decrease, reaching an almost
442 stable value (i.e., a value lower than 10% of the maximum) at
443 the eigenvalue no. 8. Similar results have been obtained for
444 other trajectories, in which a stable value is reached for
445 eigenvalues between no. 7 and no. 10. On average, the first 10
446 eigenvalues capture about 50% of the fluctuations for DBD and
447 about 40% for DBD/p28. Additionally, the first eigenvalue
448 accounts for a slightly higher percentage for DBD (23%) with
449 respect to DBD/p28 (20%). These results indicate that the
450 most relevant motions are concentrated in a few modes, whose
451 dimensionality is represented by the number of eigenvalues
452 accounting for the most relevant modes, with such an effect
453 being more relevant for DBD with respect to DBD/p28.

454 Figure 9B shows a 2D plot of the projection on the first
455 eigenvalue (PC1) versus the second one (PC2) for the two
456 trajectories of DBD (blue and magenta spots) and DBD/p28
457 (red and cyan spots); such a plot provides a global description
458 of the configurational space sampled by the system. In all of
459 the cases, the spots cover a rather wide region indicating the
460 establishment of extended collective modes in DBD; these
461 modes can be ascribed to the exploration of different structural
462 conformations. From Figure 9B, we also note that the
463 extension of the plot region is significantly reduced for
464 DBD/p28 in comparison to that of DBD alone. This means
465 that the collective modes of DBD complexed with p28 are
466 characterized by lower amplitudes, and p28 affects the global
467 dynamics of DBD by restringing the explored conformation
468 states.

469 Figure 10 shows the snapshots of the extreme projections on
470 the first eigenvectors for DBD and DBD/p28, from the same
471 two runs of Figure 9B. From this figure, it is evident that along
472 the first eigenvector, the molecule can assume different
473 arrangements throughout the whole structure, with this
474 supporting the establishment of anharmonic modes. The
475 most marked variability is observed in correspondence of the
476 H₁ and H₂ helices and of almost all turns and bends, while
477 small variability occurs at the β -sheets. Similar results have
478 been obtained for the other runs. Generally, such a behavior
479 finds correspondence with the disordered character of DBD.
480 Indeed, regions characterized by wide anharmonic modes are
481 expected to have higher flexibility, which could also confer an
482 increased capability to bind different ligands. On the other
483 hand, regions providing the skeleton of the molecule, such as
484 the β -sheets, show restricted anharmonic modes. From Figure
485 10, we also note that the extreme projections on the first
486 eigenvector are significantly reduced for DBD/p28 with
487 respect to DBD (Figure 10C,D), with this being particularly

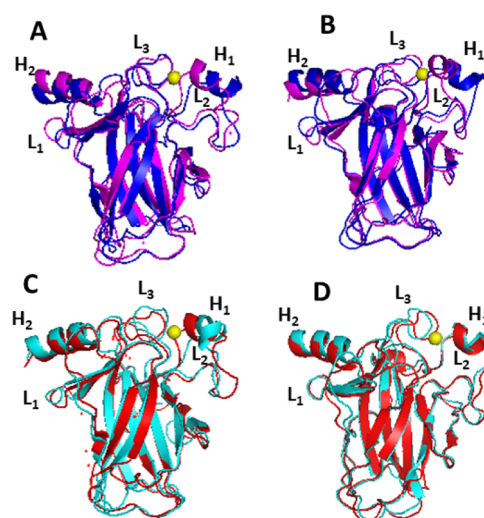


Figure 10. Extreme trajectory conformations projected along the first eigenvector for DBD (A, B) and DBD/p28 (C, D). The Zn ion is represented as a yellow sphere.

evident for the H₁ and H₂ helices, as well as at the turn located
at the top of the molecule, which is devoted to the binding of
DNA. The observed decrease in the amplitude of anharmonic
modes on DBD/p28 can be consistent with the stabilization of
the protein, which could also lead to some changes in the
binding properties of DBD. In other words, the lower
structural heterogeneity occurring in DBD/p28 with respect
to DBD could give rise to the modulation of the binding
capability of DBD, and it deserves further investigations, even
in the perspective of designing new drugs acting on the p53
functionality. Finally, it is interesting to note that the binding
of small molecules in a well-defined region of DBD can affect
the whole DBD structure, consistent with the complex
character of proteins, in which a small change might result in
drastic effects in the global system.

In summary, the ED analysis of MD trajectories shows that
the binding of p28 to DBD induces global changes in the
dynamics of DBD, yielding the reduction of collective motions.
These results can also provide a picture to interpret the
observed modulation of the heterogeneity around Trp146
when DBD is bound to p28, as emerging from fluorescence
emission data. More generally, the reduction of amplitude
collective motions could have some implications on the
capability of DBD to bind to possible partners through a global
reorganization. These aspects should be further investigated
even in connection with the possibility of enhancing the
anticancer role of p53.

CONCLUSIONS

The analysis of the time-resolved emission fluorescence of the
lone Trp146 of DBD, which constitutes the core of the tumor
suppressor p53, has put into evidence a stretched exponential
decay, indicative of a local heterogeneity around the Trp
residue. Such a result can be ascribed to the existence of
slightly different protein conformations sampled in the
analyzed temporal window and is consistent with the partially
disordered character of DBD. The increase of the stretching
parameter when DBD is complexed with the anticancer
peptide p28 reflects a reduction of the heterogeneity in the
Trp146 environment. Such a behavior suggests the capability
of p28 to modulate the conformational space explored by

528 DBD. The ED analysis of MD simulated trajectories of DBD
529 alone and of the DBD/p28 complex has shown the
530 establishment of wide collective modes throughout the DBD
531 molecule. The amplitude of these modes is wider for those
532 regions devoted to the ligand binding, such as the DNA-
533 binding domain. The reduction of the mode amplitude in the
534 DBD/p28 complex indicates that the structural heterogeneity
535 of DBD undergoes a marked modulation upon its binding to
536 the p28 peptide. Additionally, a decrease of local heterogeneity
537 around Trp146 has also been observed. These results support
538 the hypothesis that slight local changes of DBD could globally
539 affect the whole DBD dynamics. Additionally, they provide
540 some hints to analyze and improve the role of drugs able to
541 restore or enhance the anticancer capability of p53.

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553 Notes

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